

$E_{6(6)}$ Exceptional Field Theory: Applications to Type IIB Supergravity on $\text{AdS}_5 \times S^5$

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Motivations

- Search for a unified theory: reconcile matter and gravity at the quantum level.
- Important historical example: Kaluza-Klein theory.
Electromagnetism and gravity unification by embedding in a five dimensional spacetime. Provide mechanism for dimensional reduction.
- Idea still influential today $\longrightarrow$ consistent Kaluza-Klein truncation.
Powerful tool: any solution of the lower-dimensional theory can be uplifted to the higher-dimensional theory.

A toy model

Free massless scalar field in $(D+1)$ dimensional spacetime satisfy

$$\square \hat{\Phi}(x, y) = 0$$

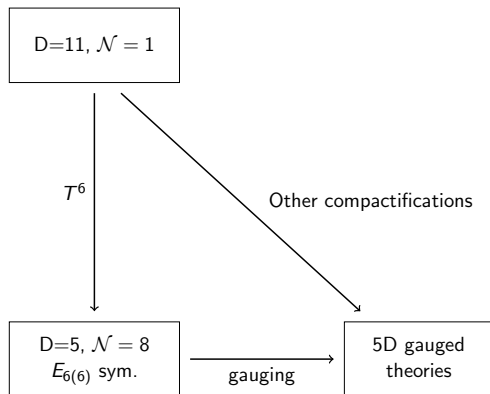
Now, suppose we compactify on circle of 'small' radius R

$$\hat{\Phi}(x, y) = \sum \phi_n(x) e^{iny/R} \rightarrow \square \phi_n(x) - \frac{n^2}{R^2} \phi_n(x) = 0$$

In this case, truncating to the massless sector ($n=0$) is consistent.
General case could be more complicated, e.g. $\square \phi_1 = \phi_0$, inconsistent.

Q: Are the fields I keep sources for the field I discard ?

Compactifications & symmetries



Dimensional reduction
→ enhanced symmetries.

Cremmer-Julia [1979]:
torus reduction of
D=11 SUGRA

$$D = 5 : E_{6(6)}$$

$$D = 4 : E_{7(7)}$$

$$D = 3 : E_{8(8)}$$

Q: Framework that make these exceptional symmetries manifest ?

The idea

Coordinate (x^μ, Y^M) , $\mu = 0..4$, M: fundamental **27** of $E_{6(6)}$

$$11 = 5 + 6$$

↓ enlarged internal space

$$5 + 27$$

Then, get rid of any redundancy with the section constraint

$$d^{MNK} \partial_N \partial_K A = 0, \quad d^{MNK} \partial_N A \partial_K B = 0$$

for any fields or gauge parameters A,B
and with d^{MNK} invariant tensor under **27**.

Content

Gauge parameter $\Lambda^M(x, Y)$: Local gauge transformations w.r.t to x-space.
 $\rightarrow \partial_\mu$ need to be covariantized.

Introduce gauge connection $\mathcal{A}_\mu^M$ and define (external) covariant derivative

$$\mathcal{D}_\mu \equiv \partial_\mu - \mathbb{L}_{\mathcal{A}_\mu}, \quad \text{such that } \delta_\Lambda \mathcal{A}_\mu^M = \mathcal{D}_\mu \Lambda^M$$

Naive curvature: $F_{\mu\nu}^M = 2\partial_{[\mu} \mathcal{A}_{\nu]}^M - [\mathcal{A}_\mu, \mathcal{A}_\nu]_E^M \rightarrow$ not covariant.

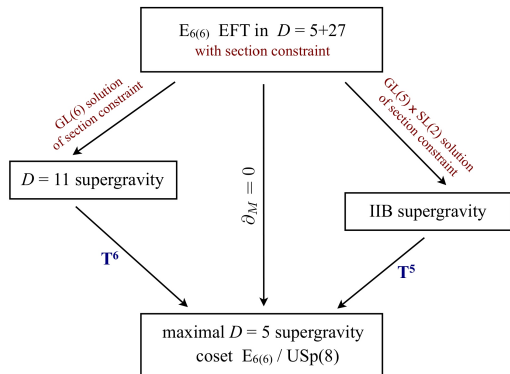
Tensor hierarchy: introduction of a two-form $\mathcal{B}_{\mu\nu M}$

$$\mathcal{F}_{\mu\nu}^M = F_{\mu\nu}^M + 10d^{MNK} \partial_K \mathcal{B}_{\mu\nu N}$$

Satisfy generalized Bianchi identity

$$3\mathcal{D}_{[\mu} \mathcal{F}_{\nu\rho]}^M = 10d^{MNK} \partial_K \mathcal{H}_{\mu\nu\rho N}$$

Solutions of the section constraints



$$E_{6(6)} \longrightarrow GL(5) \times SL(2)$$

$$27 \longrightarrow (\mathbf{5}, \mathbf{1}) \oplus (\mathbf{5}', \mathbf{2}) \oplus (\mathbf{10}, \mathbf{1}) \oplus (\mathbf{1}, \mathbf{2})$$

$$\mathcal{A}_\mu{}^M : \{ \mathcal{A}_\mu{}^m, \mathcal{A}_{\mu m \alpha}, \mathcal{A}_{\mu kmn}, \mathcal{A}_{\mu \alpha} \}$$

$$\mathcal{B}_{\mu\nu M} : \{ \mathcal{B}_{\mu\nu m}, \mathcal{B}_{\mu\nu}{}^{m\alpha}, \mathcal{B}_{\mu\nu mn}, \mathcal{B}_{\mu\nu}{}^\alpha \}$$

5+5 Kaluza-Klein decomposition

Schematically: $\{x^{\hat{\mu}}\} \rightarrow \{x^{\mu}, y^m\}$, $\hat{\mu} = 0..9$, $\mu = 0..4$, $m=1..5$

$$\begin{aligned}\hat{C}_{\hat{\mu}\hat{\nu}}^{\alpha} &\longrightarrow \{C_{\mu\nu}^{\alpha}, C_{\mu m}^{\alpha}, C_{mn}^{\alpha}\} \\ \hat{C}_{\hat{\mu}\hat{\nu}\hat{\rho}\hat{\sigma}} &\longrightarrow \{C_{\mu\nu\sigma\rho}, C_{\mu\nu\rho m}, C_{\mu\nu mn}, C_{\mu nkl}, C_{mnkl}\}\end{aligned}$$

Remaining fields after EFT fields after dualization:

$$\{\Phi, m_{mn}, b_{mn}^{\alpha}, c_{mnkl}, \mathcal{A}_{\mu m\alpha}, \mathcal{A}_{\mu kmn}, \mathcal{B}_{\mu\nu mn}, \mathcal{B}_{\mu\nu}^{\alpha}\}$$

Gauge transformation matching $\rightarrow$ IIB $\leftrightarrow$ $E_{6(6)}$ dictionary

$$\begin{aligned}A_{\mu}^m &= \mathcal{A}_{\mu}^m, \quad C_{\mu m}^{\alpha} = -\epsilon^{\alpha\beta} \mathcal{A}_{\mu m\beta}, \quad C_{\mu\nu}^{\alpha} = \tilde{\mathcal{B}}_{\mu\nu}^{\alpha}, \\ C_{\mu\nu mn} &= \frac{\sqrt{2}}{4} \tilde{\mathcal{B}}_{\mu\nu mn}, \quad C_{\mu kmn} = \frac{\sqrt{2}}{4} \mathcal{A}_{\mu kmn}\end{aligned}$$

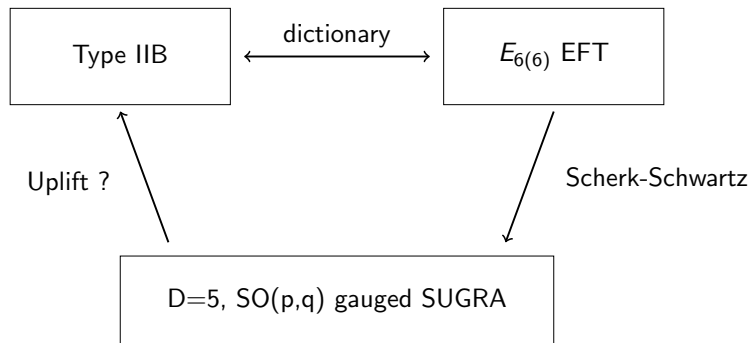
$$C_{\mu kmn} = \frac{\sqrt{2}}{4} \mathcal{Z}_{[ab] kmn}(y) A_{\mu}^{ab},$$

$$C_{\mu\nu mn} = \frac{\sqrt{2}}{4} \mathcal{K}_{[ab]}^k(y) \mathcal{Z}_{[cd] kmn}(y) A_{[\mu}^{ab}(x) A_{\nu]}^{cd}(x),$$

$$C_{m\mu\nu\rho} = -\frac{1}{32} \mathcal{K}_{[ab] m}(y) \left(2\sqrt{|g|} \varepsilon_{\mu\nu\rho\sigma\tau} M_{ab,N} F^{\sigma\tau N} + \sqrt{2} \varepsilon_{abcdef} \Omega_{\mu\nu\rho}^{cdef} \right) \\ - \frac{1}{4} \sqrt{2} \mathcal{K}_{[ab]}^k(y) \mathcal{K}_{[cd]}^l(y) \mathcal{Z}_{[ef] mkl}(y) \left(A_{[\mu}^{ab} A_{\nu}^{cd} A_{\rho]}^{ef} \right),$$

$$C_{\mu\nu\rho\sigma} = -\frac{1}{16} \mathcal{Y}_a(y) \mathcal{Y}^b(y) \left(\sqrt{|g|} \varepsilon_{\mu\nu\rho\sigma\tau} D^{\tau} M_{bc,N} M^{Nca} \right. \\ \left. + 2\sqrt{2} \varepsilon_{cdefgb} F_{[\mu\nu}^{cd} A_{\rho}^{ef} A_{\sigma]}^{ga} \right) \\ + \frac{1}{4} \left(\sqrt{2} \mathcal{K}_{[ab]}^k(y) \mathcal{K}_{[cd]}^l(y) \mathcal{K}_{[ef]}^n(y) \mathcal{Z}_{[gh] kln}(y) \right. \\ \left. - \mathcal{Y}_h(y) \mathcal{Y}^j(y) \varepsilon_{abcegj} \eta_{df} \right) A_{[\mu}^{ab} A_{\nu}^{cd} A_{\rho}^{ef} A_{\sigma]}^{gh} + \Lambda_{\mu\nu\rho\sigma}(x).$$

Conclusion



Complete proof of Kaluza-Klein consistency of $AdS_5 \times S^5$:
Any solution of the D=5 maximal $SO(p,q)$ ($p+q=6$) gauged supegravities
lift to solutions of the type IIB field equations.